

Written exam of Condensed Matter Physics - September 10th 2019
Prof. S. Caprara and A. Polimeni

Exercise 1: X ray scattering.

Sodium chloride (NaCl) crystallizes in a face-centered cubic lattice with a basis consisting of a sodium ion at $\mathbf{d}_1 = (0, 0, 0)$ and a chlorine ion at the center of the conventional cubic cell $\mathbf{d}_2 = (a/2)(1, 1, 1)$, with $a = 0.56402$ nm (see Fig. 1, left panel).

1. Determine the first 8 angles, measured with respect to the incident beam direction (see Fig. 1, right panel), for which diffraction peaks are observed on a detector using the powder or Debye-Scherrer method. The x-ray beam wavelength is $\lambda = 0.15$ nm.
2. Say which peaks are more intense assuming that the atomic form factor f can be put equal to Z (i.e., the atomic number) times the amplitude of the wave scattered from one electron.

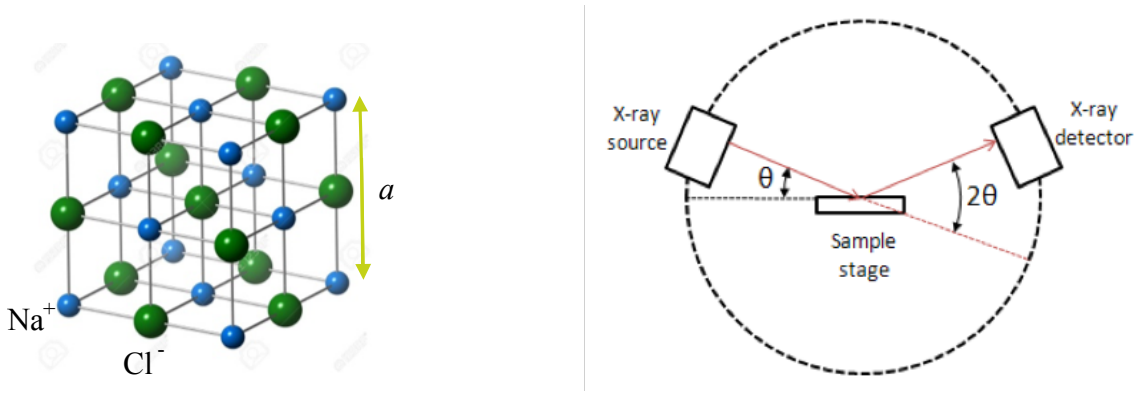


Fig. 1.

Exercise 2: Intrinsic semiconductors.

Consider a two-dimensional intrinsic semiconductor, that is described by the following density of states

$$D(\varepsilon) = \begin{cases} D_v, & \text{for } \varepsilon < \varepsilon_v, \\ D_c, & \text{for } \varepsilon > \varepsilon_c, \\ 0, & \text{elsewhere,} \end{cases}$$

such that the number of states per unit surface, in the interval $[\varepsilon, \varepsilon + d\varepsilon]$, is $D(\varepsilon) d\varepsilon$, with constant D_v and D_c , ε_v and ε_c being the thresholds of the valence and conduction band, respectively.

1. Determine the density of electrons in the conduction band, n , and the density of holes in the valence band, p , as a function of $D_v, D_c, \varepsilon_v, \varepsilon_c$, at a generic temperature T (assume, here and in the following, that the temperature is such that $\mu - \varepsilon_v, \varepsilon_c - \mu \gg \kappa_B T$, where μ is the chemical potential and κ_B is the Boltzmann constant).
2. Determine the expression of the chemical potential μ , as a function of $D_v, D_c, \varepsilon_v, \varepsilon_c$, at a generic temperature T .
3. Let now $D_v = 2.5 \times 10^{18} \text{ eV}^{-1} \cdot \text{m}^{-2}$, $D_c = 5.0 \times 10^{18} \text{ eV}^{-1} \cdot \text{m}^{-2}$, $\varepsilon_v = 1.05 \text{ eV}$, $\varepsilon_c = 1.45 \text{ eV}$, evaluate n, p and μ at $T = 300 \text{ K}$.

[Note that 1 eV corresponds to an energy of $1.60 \times 10^{-19} \text{ J}$; the Boltzmann constant is $\kappa_B = 1.38 \times 10^{-23} \text{ J} \cdot \text{K}^{-1}$].

Solution of the written exam
Prof. S. Caprara and A. Polimeni

Exercise 1.

1. Since

$$\theta = \arcsin\left(\frac{\lambda}{2d}\right) = \arcsin\left(\frac{\lambda}{a}\sqrt{h^2 + k^2 + l^2}\right),$$

where the indices h, k, l are all even or all odd, we find

	h	k	l	d (nm)	θ (deg)	2θ (deg)
	1	1	1	0.32564	13.31578	26.63156
	2	2	0	0.28201	15.42329	30.84658
	3	2	2	0.19941	22.09277	44.18553
	4	3	1	0.17006	26.16926	52.33858
	5	2	2	0.16282	27.42812	54.85624
	6	4	0	0.14142	32.13365	64.26730
	7	3	3	0.12940	35.42384	70.84768
	8	4	2	0.12612	36.48967	72.97934

2. The resulting scattered amplitude is

$$\begin{aligned} I &= A \left| f_{\text{Na}} \left[e^{-i0} + e^{-i\pi(h+k)} + e^{-i\pi(h+l)} + e^{-i\pi(k+l)} \right] \right. \\ &\quad \left. + f_{\text{Cl}} \left[e^{-i\pi(h+k+l)} + e^{-i\pi(2h+2k+l)} + e^{-i\pi(2h+k+2l)} + e^{-i\pi(h+2k+2l)} \right] \right|^2 \\ &= A \left| \left[f_{\text{Na}} + f_{\text{Cl}} e^{-i\pi(h+k+l)} \right] \left[1 + e^{-i\pi(h+k)} + e^{-i\pi(h+l)} + e^{-i\pi(k+l)} \right] \right|^2. \end{aligned}$$

The scattered amplitude is non zero provided the h, k, l indices are all even, and $I = 16 A (f_{\text{Na}} + f_{\text{Cl}})^2$ (greater intensity), or all odd, and $I = 16 A (f_{\text{Na}} - f_{\text{Cl}})^2$ (lower intensity).

Exercise 2.

1. We have

$$n = \int_{\varepsilon_c}^{+\infty} \frac{D_c}{e^{(\varepsilon-\mu)/\kappa_B T} + 1} d\varepsilon \approx D_c \int_{\varepsilon_c}^{+\infty} e^{-(\varepsilon-\mu)/\kappa_B T} d\varepsilon = \kappa_B T D_c e^{-(\varepsilon_c-\mu)/\kappa_B T}, \quad (1)$$

and

$$p = \int_{-\infty}^{\varepsilon_v} D_v \left[1 - \frac{1}{e^{(\varepsilon-\mu)/\kappa_B T} + 1} \right] d\varepsilon \approx D_v \int_{-\infty}^{\varepsilon_v} e^{(\varepsilon-\mu)/\kappa_B T} d\varepsilon = \kappa_B T D_v e^{(\varepsilon_v-\mu)/\kappa_B T}. \quad (2)$$

Multiplying Eq. (1) by Eq. (2), the dependence on μ drops and we find

$$np = (\kappa_B T)^2 D_c D_v e^{(\varepsilon_v-\varepsilon_c)/\kappa_B T} \equiv n_i^2,$$

which is the *law of mass action* in the present situation. In an intrinsic semiconductor

$$n = p = n_i = \kappa_B T \sqrt{D_c D_v} e^{(\varepsilon_v-\varepsilon_c)/2\kappa_B T}.$$

2. Dividing Eq. (1) by Eq. (2), and taking into account that $n = p$ in an intrinsic semiconductor, we find

$$1 = \frac{D_c}{D_v} e^{(2\mu-\varepsilon_v-\varepsilon_c)/\kappa_B T} \quad \Rightarrow \quad 0 = \log \frac{D_c}{D_v} + \frac{2\mu - \varepsilon_v - \varepsilon_c}{\kappa_B T},$$

hence

$$\mu = \frac{1}{2} \left(\varepsilon_v + \varepsilon_c + \kappa_B T \log \frac{D_v}{D_c} \right).$$

3. At $T = 300$ K, we have $\kappa_B T = 0.0259$ eV, hence

$$p = n = n_i = 4.06 \times 10^{13} \text{ m}^{-2},$$

and

$$\mu = 1.24 \text{ eV}.$$