Written exam of Condensed Matter Physics - June 16th 2021 Profs. S. Caprara and A. Polimeni

Exercise 1: Phonons [15 points].

A metal with valence 3 features a simple cubic cell with a single-atom basis and an electron density $n_{\rm el} = 3.66 \times 10^{29} \,\mathrm{m}^{-3}$. The electrons can be treated as a gas of free particles. Concerning the metal ion vibrations, we have three acoustic branches (one longitudinal L and two transverse T) along the direction of the reciprocal lattice side:

$$\omega_{\rm T} = \omega_{\rm T}^0 \sin\left(\frac{qa}{2}\right), \qquad \omega_{\rm L} = \omega_{\rm L}^0 \sin\left(\frac{qa}{2}\right),$$

with $\omega_{\rm L}^0 = 4.8 \times 10^{12} \, {\rm rad/s}$. The Debye temperature of the transverse mode is $\Theta_{\rm T}^{\rm D} = 45.7 \, {\rm K}$.

1. Determine the lattice constant a and the Debye wavevector $q_{\rm D}$ [5 points].

2. Determine the velocity of sound for longitudinal and transverse modes [5 points].

3. The low temperature specific heat at constant volume is given by $c_V(T) = AT + BT^3$.

Determine the value of A, B and of c_V at a temperature T = 1 K, knowing that the Fermi energy of the metal is $E_F = 18.64 \text{ eV}$ [5 points].



Fig. 1.

Exercise 2: Tight binding [15 points].

A two-dimensional rectangular crystal with lattice parameters a = 0.4 nm and b = 0.3 nm hosts the compound with chemical formula QR. The position of the Q atom within the primitive cell is identified by the basis vector $d_{\rm Q} = (0,0)$, while the position of the R atom is identified by the basis vector $d_{\rm R} = (\frac{a}{2}, \frac{b}{2})$. Assume that the electron states can be described within the tight binding model with attractive potential $\Delta U < 0$ and both atoms contribute to the formation of the relevant electron bands with s-type orbitals (see Fig. 1). The only transfer integral to be considered is $\gamma = 0.8 \,\mathrm{eV}$, between Q atoms and nearest-neighboring R atoms (and viceversa). All other transfer integrals and all overlap integrals can be neglected. For simplicity, put to zero all the β integrals, and take the atomic levels $\varepsilon_{\rm Q} = 1.0 \,\mathrm{eV}$ and $\varepsilon_{\rm R} = 3.0 \,\mathrm{eV}$.

1. Determine the energy vs. quasi-momentum dispersion relations $E_{\pm}(\mathbf{k})$, where + labels the conduction (upper) band and - labels the valence (lower) band, and $\mathbf{k} = (k_x, k_y)$ [5 points].

2. Determine the position in quasi-momentum space and the energy of the maximum of the valence band and of the minimum of the conduction band. Determine the energy gap $E_{\rm g}$ between the valence band and the conduction band [3 points].

3. Determine the expression and numerical values of the elements of the effective mass tensors for the two bands at the $\Gamma = (0,0)$ point of the Brillouin zone [7 points].

[Note that 1 eV corresponds to an energy of 1.602×10^{-19} J; the Planck constant is $\hbar = 1.055 \times 10^{-34}$ J·s; the Boltzmann constant is $\kappa_B = 1.381 \times 10^{-23}$ J·K⁻¹; the free electron mass is $m_0 = 9.109 \times 10^{-31}$ kg].

Solution of the written exam Profs. S. Caprara and A. Polimeni

Exercise 1.

1. The lattice constant is given by $a = n_{\rm at}^{-1/3} = 0.2016 \,\mathrm{nm}$, where the atomic density is $n_{\rm at} = n_{\rm el}/3$. The Debye wavevector is $q_{\rm D} = (6\pi^2 n_{\rm at})^{1/3} = 1.933 \times 10^{10} \,\mathrm{m}^{-1}$.

2. The longitudinal sound velocity can be estimated as

$$v_{\rm L} = \lim_{q \to 0} \frac{\mathrm{d}\omega_{\rm L}}{\mathrm{d}q} = \frac{1}{2} a \omega_{\rm L}^0 = 483.8 \,\mathrm{m/s}.$$

The transverse sound velocity is

$$v_{\mathrm{T}} = \frac{\kappa_B \Theta_{\mathrm{T}}^{\mathrm{D}}}{\hbar q_{\mathrm{D}}} = 309.5 \,\mathrm{m/s}.$$

3. We know that the lattice contribution is BT^3 , where

$$B = \frac{2\pi^2}{15} \frac{\kappa_B^4}{\hbar^3} \left(\frac{1}{v_{\rm L}^3} + \frac{2}{v_{\rm T}^3} \right) = 3.110 \times 10^3 \, \frac{\rm J}{\rm m^3 \cdot \rm K^4},$$

while the electronic contribution is AT, where

$$A=\frac{\pi^2}{2}\frac{\kappa_B^2}{E_{\rm F}}n_{\rm el}=115.3\,\frac{\rm J}{\rm m^3\cdot K^2}$$

Therefore, the specific heat at T = 1 K is $c_V = 3.225 \times 10^3 \text{ J/(m^3 \text{ K})}$.

Exercise 2.

1. The four vectors that locate the nearest-neighbor sites in the given lattice are $\mathbf{R} = (\pm \frac{a}{2}, \pm \frac{b}{2})$. Let

$$g_{\mathbf{k}} \equiv \gamma \sum_{\mathbf{R}} e^{i\mathbf{k}\cdot\mathbf{R}} = 4\gamma \cos\left(\frac{ak_x}{2}\right) \cos\left(\frac{bk_y}{2}\right).$$

Then, the coefficients $b_{\rm Q}, b_{\rm R}$ of the linear combination of atomic orbitals within the tight-binding method obey the set of linear equations

$$\begin{cases} \left[\varepsilon_{\mathrm{Q}} - E(\boldsymbol{k})\right] b_{\mathrm{Q}} - g_{\boldsymbol{k}} b_{\mathrm{R}} &= 0, \\ \left[\varepsilon_{\mathrm{R}} - E(\boldsymbol{k})\right] b_{\mathrm{R}} - g_{\boldsymbol{k}} b_{\mathrm{Q}} &= 0, \end{cases}$$

which has nontrivial solutions only if

$$E(\boldsymbol{k}) = \frac{\varepsilon_{\mathrm{Q}} + \varepsilon_{\mathrm{R}}}{2} \pm \sqrt{\left(\frac{\varepsilon_{\mathrm{R}} - \varepsilon_{\mathrm{Q}}}{2}\right)^2 + g_{\boldsymbol{k}}^2} \equiv E_{\pm}(\boldsymbol{k}).$$

2. The location of the maxima (minima) of the valence (conduction) band is given by the condition $g_{\mathbf{k}} = 0$, i.e., $k_x = \pm \frac{\pi}{a}$ or $k_y = \pm \frac{\pi}{b}$. These conditions correspond to the boundary of the first Brillouin zone. The values of the band energies are $E_+ = \varepsilon_{\rm R} = 3.0 \,\mathrm{eV}$ and $E_- = \varepsilon_{\rm Q} = 1.0 \,\mathrm{eV}$. The gap is $E_{\rm g} = \varepsilon_{\rm R} - \varepsilon_{\rm Q} = 2.0 \,\mathrm{eV}$.

3. Near the Γ point $g_{\boldsymbol{k}} \approx 4\gamma \left(1 - \frac{a^2 k_x^2 + b^2 k_y^2}{8}\right)$, and $g_{\boldsymbol{k}}^2 \approx 16\gamma^2 \left(1 - \frac{a^2 k_x^2 + b^2 k_y^2}{4}\right)$. Then

$$E_{\pm} \approx \frac{\varepsilon_{\mathrm{R}} + \varepsilon_{\mathrm{Q}}}{2} \pm \sqrt{\left(\frac{\varepsilon_{\mathrm{R}} - \varepsilon_{\mathrm{Q}}}{2}\right)^2 + 16\gamma^2 - 4\gamma^2 (a^2 k_x^2 + b^2 k_y^2)}$$
$$\approx \frac{\varepsilon_{\mathrm{R}} + \varepsilon_{\mathrm{Q}}}{2} \pm \sqrt{\left(\frac{\varepsilon_{\mathrm{R}} - \varepsilon_{\mathrm{Q}}}{2}\right)^2 + 16\gamma^2} \left[1 - \frac{2\gamma^2 (a^2 k_x^2 + b^2 k_y^2)}{\left(\frac{\varepsilon_{\mathrm{R}} - \varepsilon_{\mathrm{Q}}}{2}\right)^2 + 16\gamma^2}\right].$$

Focusing on the k-dependent part, we have

$$E_{\pm} \approx \ldots \mp \frac{2\gamma^2 (a^2 k_x^2 + b^2 k_y^2)}{\sqrt{\left(\frac{\varepsilon_{\mathrm{R}} - \varepsilon_{\mathrm{Q}}}{2}\right)^2 + 16\gamma^2}} = \ldots \mp \left(\frac{\hbar^2 k_x^2}{2m_{xx}^*} + \frac{\hbar^2 k_y^2}{2m_{yy}^*}\right),$$

so the mass tensor at the Γ point is diagonal, with equal absolute value of the effective masses along the principal axes for both bands,

$$m_{xx}^{*} = \frac{\hbar^{2} \sqrt{\left(\frac{\varepsilon_{\mathrm{R}} - \varepsilon_{\mathrm{Q}}}{2}\right)^{2} + 16\gamma^{2}}}{4\gamma^{2}a^{2}} = \sqrt{1 + \left(\frac{\varepsilon_{\mathrm{R}} - \varepsilon_{\mathrm{Q}}}{8\gamma}\right)^{2}} \frac{\hbar^{2}}{\gamma a^{2}} = \frac{\sqrt{281}}{16} \frac{\hbar^{2}}{\gamma a^{2}} = 5.68 \times 10^{-31} \,\mathrm{kg} = 0.624 \,m_{0},$$

and

$$m_{yy}^{*} = \frac{\hbar^{2} \sqrt{\left(\frac{\varepsilon_{\mathrm{R}} - \varepsilon_{\mathrm{Q}}}{2}\right)^{2} + 16\gamma^{2}}}{4\gamma^{2}b^{2}} = \sqrt{1 + \left(\frac{\varepsilon_{\mathrm{R}} - \varepsilon_{\mathrm{Q}}}{8\gamma}\right)^{2}} \frac{\hbar^{2}}{\gamma b^{2}} = \frac{\sqrt{281}}{16} \frac{\hbar^{2}}{\gamma b^{2}} = 1.01 \times 10^{-30} \,\mathrm{kg} = 1.11 \,m_{0}$$