## Written exam of Condensed Matter Physics - November 4th 2022 Profs. S. Caprara and A. Polimeni

## Exercise 1: Phonons.

Consider a one-dimensional monoatomic lattice having length L = 40 cm. The phonon density of states, within the Debye model, is a constant  $\mathcal{D} = 8.00 \times 10^{-5}$  states/(rad/s), for  $0 \le \omega \le \omega_{\rm D}$ , where  $\omega_{\rm D}$  is the Debye frequency.

1. [4 points] Find the velocity of sound  $v_s$  of the given lattice.

2. [3 points] Evaluate the number of atoms N in the lattice, given the Debye frequency  $\omega_{\rm D} = 4.00 \times 10^{12} \, \text{rad/s}$ .

3. [4 points] Determine the mass m of the atoms forming the lattice, knowing that the elastic constant that mimics the interaction between nearest-neighboring atoms is  $\beta = 8.00 \times 10^{-3} \text{ N/m}$ .

3. [4 points] Evaluate the mean energy per unit length u at a temperature T = 500 K, after checking that  $\hbar\omega_{\rm D} \ll \kappa_{\rm B}T$ , so that the classical approximation holds.

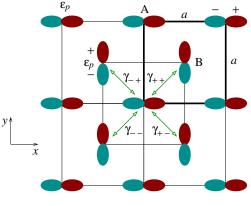


Fig. 1.

#### Exercise 2: Tight binding.

A binary compound AB forms a two-dimensional square Bravais lattice with lattice spacing a = 0.5 nm (see Fig. 1). The atoms A occupy the sites of the Bravais lattice, while the atoms B are located at the positions identified by the basis vector  $\boldsymbol{d} = (a/2, a/2)$ . A primitive cell is highlighted by a thicker line in Fig. 1. The atoms A display outer  $p_x$  orbitals, while the atoms B display outer  $p_y$  orbitals. The convention for the signs of the lobes of the p orbitals is shown in Fig. 1. The energy of both p orbitals is  $\varepsilon_p = 0$ . Assume that the electron states of the given lattice can be described within the tight-binding approximation with attractive lattice potential  $\Delta U < 0$ , and consider only nearest-neighbor transfer integrals  $\gamma_{uv} \equiv \gamma \left( \boldsymbol{r}_{uv} \equiv u \frac{a}{2} \hat{\boldsymbol{x}} + v \frac{a}{2} \hat{\boldsymbol{y}} \right)$ , where  $u = \pm$ ,  $v = \pm$ , and  $\hat{\boldsymbol{x}}$ ,  $\hat{\boldsymbol{y}}$  are the unit vectors of the corresponding axes, so that the four vectors  $\boldsymbol{r}_{uv}$  locate the positions of the nearest neighbors. The absolute value of the nearest-neighbor transfer integrals is  $\gamma_0 \equiv |\gamma_{uv}| = 0.75 \,\text{eV}$ . All other transfer integrals and all on-site ( $\beta$ ) and overlap ( $\alpha$ ) integrals can be neglected.

1. [3 points] Assign the correct sign to the four nearest-neighbor transfer integrals,  $\gamma_{++}$ ,  $\gamma_{+-}$ ,  $\gamma_{-+}$ , and  $\gamma_{--}$  (see Fig. 1), motivating your answer.

2. [5 points] Determine the two electron bands formed by the orbitals taken into account,  $\varepsilon_{\pm}(\mathbf{k})$ , where  $\mathbf{k} = (k_x, k_y)$  is the two-dimensional Bloch wave vector.

3. [7 points] Determine the location of the minima of the lowest band  $\varepsilon_{-}(\mathbf{k})$  in the first Brillouin zone, then calculate the numerical values of the eigenvalues of the effective mass tensor  $m_{ij}$ , i = x, y, j = x, y, at these points.

[Useful constants and conversion factors: the Planck constant is  $\hbar = 1.05 \times 10^{-34} \,\text{J}\cdot\text{s}$ , the Boltzmann constant is  $\kappa_B = 1.38 \times 10^{-23} \,\text{J}\cdot\text{K}^{-1}$ , the elementary charge is  $e = 1.60 \times 10^{-19} \,\text{C}$ , the free electron mass is  $m_0 = 9.11 \times 10^{-31} \,\text{kg}$ ; 1 eV corresponds to a temperature of  $1.16 \times 10^4 \,\text{K}$  or to an energy of  $1.60 \times 10^{-19} \,\text{J}$ ].

# Solution Profs. S. Caprara and A. Polimeni

Exercise 1.

1. One has

$$\mathcal{D}(\omega) = \frac{\mathrm{d}N}{\mathrm{d}\omega} = \frac{L}{\pi} \frac{\mathrm{d}q}{\mathrm{d}\omega} = \frac{L}{\pi v_s},$$

for  $0 \leq \omega \leq \omega_{\rm D}$ . Hence, one finds

$$v_s = \frac{L}{\pi \mathcal{D}} = 1.592 \times 10^3 \,\mathrm{m/s}$$

2. One has

$$N = \int_0^{\omega_{\rm D}} \mathcal{D} \,\mathrm{d}\omega = \mathcal{D} \,\omega_{\rm D} = 3.2 \times 10^8.$$

3. One has  $v_s = a\sqrt{\beta/m}$ , where a = L/N. Hence,

$$m = \left(\frac{L}{N}\right)^2 \frac{\beta}{v_s^2} = 4.93 \times 10^{-27} \,\mathrm{kg}$$

4. One finds  $\hbar\omega_{\rm D} = 4.22 \times 10^{-22} \,\text{J} \ll \kappa_{\rm B}T = 6.90 \times 10^{-21} \,\text{J}$ , hence the classical approximation can be used and

$$u = \frac{\kappa_{\rm B}T}{a} = \frac{N\kappa_{\rm B}T}{L} = 5.52 \times 10^{-12} \,\mathrm{J/m}$$

## Exercise 2.

1. From Fig. 1 one can see that the positive lobe of the  $p_x$  orbital of an A atom is closer to the positive lobe of the  $p_y$  orbital a B atom at  $\mathbf{r}_{+-}$  and to the negative lobe of the  $p_y$  orbital a B atom at  $\mathbf{r}_{++}$ , whereas the negative lobe of the  $p_y$  orbital of an A atom is closer to the negative lobe of the  $p_y$  orbital a B atom at  $\mathbf{r}_{-+}$  and to the positive lobe of the  $p_y$  orbital a B atom at  $\mathbf{r}_{-+}$  and to the positive lobe of the  $p_y$  orbital a B atom at  $\mathbf{r}_{-+}$  and to the positive lobe of the  $p_y$  orbital a B atom at  $\mathbf{r}_{-+}$  and to the positive lobe of the  $p_y$  orbital a B atom at  $\mathbf{r}_{--}$ .

2. Let

$$w_{\boldsymbol{k}} = \sum_{u=\pm} \sum_{v=\pm} \gamma_{u,v} \mathrm{e}^{\mathrm{i}\boldsymbol{r}_{u,v}\cdot\boldsymbol{k}} = 2\gamma_0 \left\{ \cos\left[\frac{a(k_x - k_y)}{2}\right] - \cos\left[\frac{a(k_x + k_y)}{2}\right] \right\} = 4\gamma_0 \sin\left(\frac{k_x a}{2}\right) \sin\left(\frac{k_y a}{2}\right).$$

Indicating with  $b_x$  and  $b_y$  the coefficients of the linear combination of  $p_x$  and  $p_y$  orbitals, respectively, within the tight-binding approximation indicated in the text, one find the text of two coupled homogeneous linear equations

$$\begin{cases} (\varepsilon_p - \varepsilon_k)b_x - w_k \, b_y = 0, \\ -w_k b_x + (\varepsilon_p - \varepsilon_k)b_x = 0, \end{cases}$$

which has nontrivial solutions if  $\varepsilon_{\mathbf{k}} \equiv \varepsilon_{\pm}(\mathbf{k}) = \varepsilon_p \pm |w_{\mathbf{k}}|$ . In the following, according to the text, we set  $\varepsilon_p = 0$ . 3. The minima of the lower band,  $\varepsilon_{-}(\mathbf{k})$ , are found where

$$\left|\sin\left(\frac{k_x a}{2}\right)\sin\left(\frac{k_y a}{2}\right)\right|$$

is maximum. In the first Brillouin zone  $k_x \in [-\frac{\pi}{a}, \frac{\pi}{a}], k_y \in [-\frac{\pi}{a}, \frac{\pi}{a}]$ , this occurs at the four equivalent points  $\mathbf{k}_{\pm\pm} = (\pm \frac{\pi}{a}, \pm \frac{\pi}{a})$ , where the two sines are both equal to  $\pm 1$ . Near the minima, then,

$$\varepsilon_{-}(\mathbf{k}) \approx -4\gamma_{0} + \frac{1}{2}a^{2}\gamma_{0}\left(\delta k_{x}^{2} + \delta k_{y}^{2}\right),$$

where  $\delta k_x$  and  $\delta k_y$  are the deviations of  $k_x$  and  $k_y$  with respect to their values at a given minimum. Hence, the inverse mass tensor is diagonal, and the two eigenvalues are equal, according to the symmetries of a square lattice. Thus, one finds

$$m_{xx} = m_{yy} = \frac{\hbar^2}{a^2 \gamma_0} = 3.70 \times 10^{-31} \,\mathrm{kg} = 0.406 \,m_0.$$