

IMPORTANCE OF ELECTRON-ELECTRON INTERACTION

EXAMPLE	IONIZATION ENERGY OF NOBLE GASES	LAST OCCUPIED SHELL	IONIZATION ENERGY - INDEX (eV) ELECTRON $\frac{E^2}{n^2}$	EXP. IONIZ. (eV)
He	2	$n=1$	4	1.81
Ne	10	$n=2$	25	1.59
Ar	18	$n=3$	36	1.16
Kr	36	$n=4$	64	1.03
Xe	54	$n=5$	116, 64	0.89

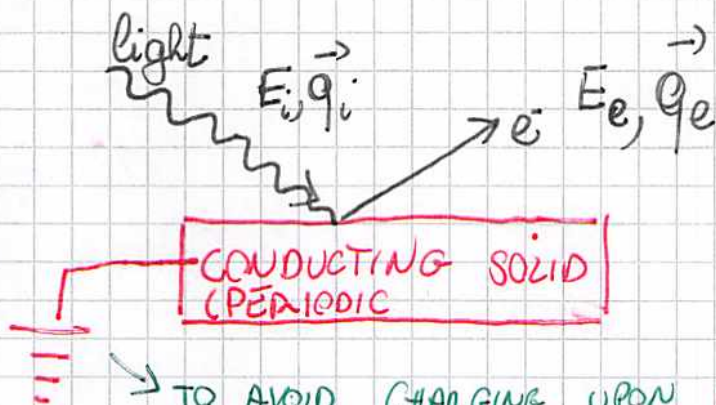
ELECTRONIC BANDS BASED ON BLOCH-THEOREM

VALID IN A SINGLE PARTICLE PICTURE

DO THE ELECTRONIC BAND EXISTS IN REALITY?

ANGLE RESOLVED PHOTOEMISSION SPECTROSCOPY (ARPES)

(IONISATION OF A SOLID)



→ TO AVOID CHARGING UPON MULTIPLE IONIZATION AND (COULOMB INTERACTION BETWEEN ESCAPING ELECTRON AND CHARGED SOLID)

ELECTRON - ESCAPES ONLY FROM FEW ATOMIC PLANE SURFACE OF THE SOLID PERIODIC OR 2D MATERIAL (ONLY IN PLANE PERIODICITY) $\left\{ \vec{G}_{\parallel} \right\}$ RECIPROCAL LATTICE OF THE SURFACE

- CONSERVATION OF ENERGY (GOOD QUANTUM NUMBER)

IN THE MATERIAL WE DEPOSIT AN ENERGY ^{GAS DISCHARGE}

$$E = E_i - E_e$$

$$E_i = 20-50 \text{ eV (GAS DISCHARGE)}$$

$$10-100 \text{ KeV (SYNCHROTRON)}$$

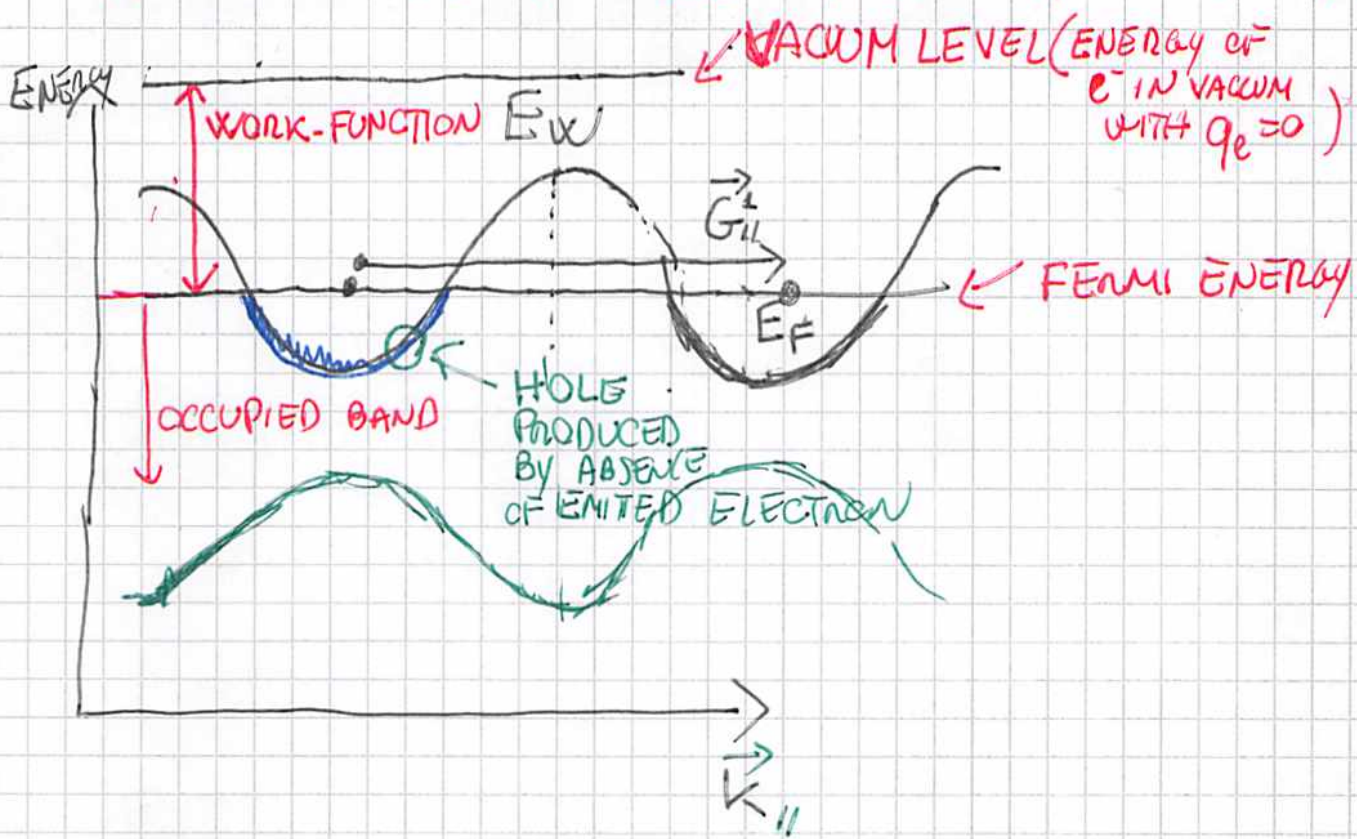
$$+ 6-7 \text{ eV (LASER)}$$

- A QUASIMOMENTUM (

$$\vec{q}_{\parallel} = \vec{q}_i - \vec{q}_e + \vec{G}_{\parallel}$$

- WE REMOVE ONE ELECTRON

INDEPENDENT ELECTRON INTERPRETATION



EMITTED ELECTRON LEAVES A HOLE

WITH MOMENTUM $\vec{q}_{||} = \vec{q}_{||} - \vec{q}_{e||} + \vec{G}_{||}$

AND BAND ENERGY $\vec{E}_{\vec{q}_{||}} = E_e - E_i + E_w$
WORK-FUNCTION CONSTANT
 WHERE $E_F = 0$

IN THIS PICTURE WE PROBE THE OCCUPIED ENERGY BANDS

TRUE IN A INDEP-PARTICLE PICTURE

DO WE SEE BANDS? (DOES BLOCH-THEOREM WORK?)

WHY IN PRESENCE OF A
HUGE e-e INTERACTION (SIMILAR TO
THE e-NUCLEUS INTERACTION) WE
SEE BANDS?

- ELECTRONS MOVE IN AN "EFFECTIVE"
ONE BODY POTENTIAL" GENERATED BY
NUCLEI AND BY OTHER ELECTRONS
COMPARABLE SIZE

- WE CAN DEFINE IN A RIGOROUS
WAY SUCH "EFFECTIVE POTENTIAL"
AND COMPUTE IT NUMERICALLY

- DIFFERENT APPROACHES (DIFFERENT POTENTIAL)
- FOR TOTAL ENERGY HF (LOW ACCURACY)
DFT (HIGH ACCURACY)
- FOR BAND STRUCTURE DFT (MEDIUM ACCURACY)
GW (HIGH ACCURACY)

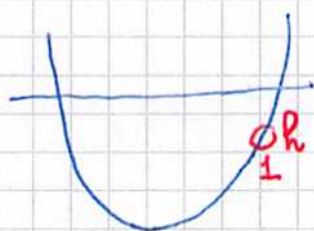
IN THE SLIDES COMPARISON BETWEEN DFT
AND ARPES (DFT FOR GRAPHENE RECALIBRATING ~15%
REQUIRED)

SIGNATURE OF ELECTRON-ELECTRON INTERACTION

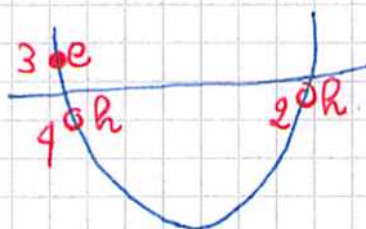
DECAY OF A HOLE (SAME FOR ELECTRON BUT ARPES CREATES HOLES!)



INITIAL STATE



FINAL STATE



CONSERVATION OF ENERGY AND QUASIMOMENTUM BY e-e INTERACTION

$$\text{ENERGY} \quad \underbrace{-E_{\vec{k}_1}}_{\text{INITIAL}} = \underbrace{-E_{\vec{k}_2} - E_{\vec{k}_4} + E_{\vec{k}_3}}_{\text{FINAL}}$$

$$\text{QUASI-MOMENTUM} \quad -\vec{k}_1 = -\vec{k}_2 - \vec{k}_4 + \vec{k}_3 + \vec{G}$$

← ARBITRARY REC LATTICE VECTOR (INCLUDING $\vec{G}=0$)

BROADENING OF BAND AT \vec{k}_1 BY ARPES

$$\frac{\Gamma_{\vec{k}_1}}{\hbar} = \frac{1}{\tau_{h_1}} = \frac{2\pi}{\hbar} \int_{\substack{\vec{k}_2 < E_F \\ \vec{k}_4 < E_F \\ \vec{k}_3 > E_F}} \frac{d^3k_2}{(2\pi)^3} \frac{d^3k_3}{(2\pi)^3} \frac{d^3k_4}{(2\pi)^3} \delta(-E_{\vec{k}_1} + E_{\vec{k}_2} + E_{\vec{k}_4} - E_{\vec{k}_3}) \times \sum_{\vec{G}} \delta(-\vec{k}_1 + \vec{k}_2 + \vec{k}_4 - \vec{k}_3 + \vec{G}) \times (\text{MATRIX ELEMENT})^2$$

↑
LIFETIME OF HOLE

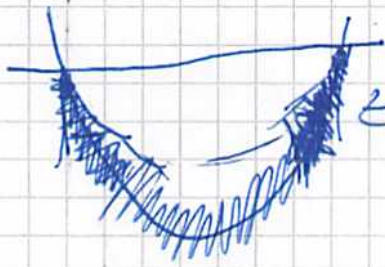
$$\left[\text{FERMI GOLDEN RULE } \Gamma_i = \frac{2\pi}{\hbar} \sum_f |\langle f | H^{\text{int}} | i \rangle|^2 \delta(E_f - E_i) \right]$$

$AT=0$ IN 3D (FOR BOTH HOLES AND ELECTRONS)

$$\Gamma_{k_1} \propto (\epsilon_F - \epsilon_{k_1})^2$$

IN 2D

$$\Gamma_{k_2} \propto (\epsilon_F - \epsilon_{k_2})^2 \ln \left(\frac{\epsilon_F}{|\epsilon_F - \epsilon_{k_2}|} \right)$$

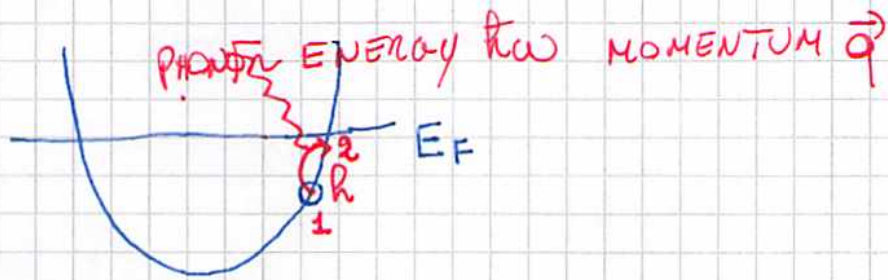


← BROADENING INCREASES WITH THE DISTANCE FROM THE FERMI ENERGY

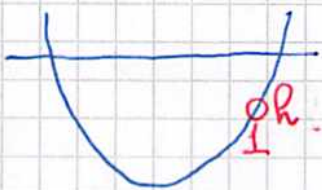
SIGNATURE OF ELECTRON-PHONON INTERACTION

DECAY OF A HOLE INTO A NON-DISPERSIVE OPTICAL PHONON (EINSTEIN-MODE)

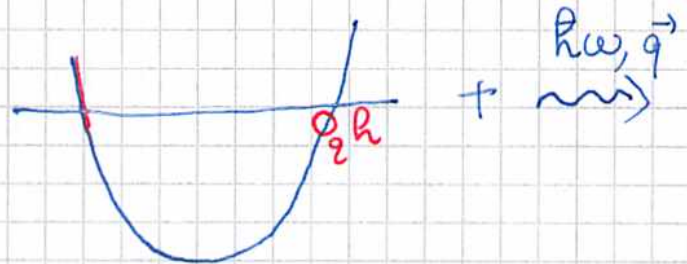
$$\hbar\omega_{\vec{q}} = \hbar\omega$$



INITIAL STATE



FINAL STATE



CONSERVATION OF ENERGY

$$-E_{\vec{k}_1} = -E_{\vec{k}_2} + \hbar\omega \quad \left| \quad \text{ONLY POSSIBLE IF } (E_F - E_{\vec{k}_1}) > \hbar\omega \right.$$

CONSERVATION OF MOMENTUM

$$-\vec{k}_1 = -\vec{k}_2 + \vec{q} + \vec{G} \quad \left. \begin{array}{l} \vec{G} \in \text{RECIPROCAL LATTICE} \\ \text{IF } \vec{G} = 0 \quad \vec{q} = \vec{k}_2 - \vec{k}_1 \end{array} \right.$$

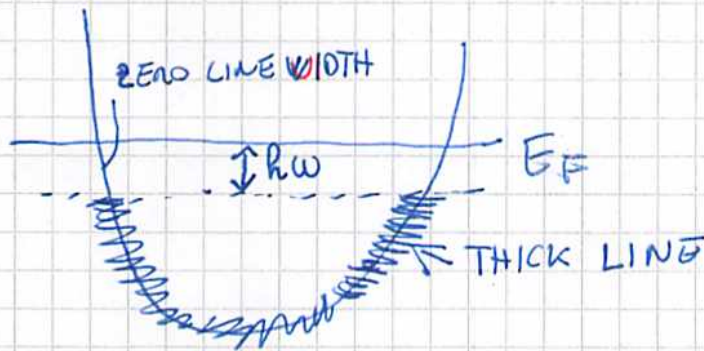
$$\frac{\Gamma_{\vec{k}_1}}{\hbar} = \frac{1}{\tau_{\vec{k}_1}} \stackrel{\text{FERMI GOLDEN RULE}}{=} \frac{2\pi}{\hbar} \int_{E_{\vec{k}_2} < E_F} \frac{d^3k_2}{(2\pi)^3} \delta[E_{\vec{k}_2} - (E_{\vec{k}_1} + \hbar\omega)] M^2$$

IF M^2 INDEPENDENT OF \vec{k}_1, \vec{k}_2

$$\Gamma_{\vec{k}_1}^2 = 2\pi M^2 \int_{\substack{\frac{d^3 k_2}{(2\pi)^3} \\ \epsilon_F - \epsilon_{\vec{k}_1} > \hbar\omega}} \delta[\epsilon_{\vec{k}_2} - (\epsilon_{\vec{k}_1} + \hbar\omega)]$$

$$= 2\pi M_x^2 \begin{cases} 0 & \text{IF } \epsilon_F - \epsilon_{k_1} < \hbar\omega \\ \frac{1}{2} \text{ DOS}(\epsilon_{k_1} + \hbar\omega) & \end{cases}$$

SPIN \swarrow \searrow
ELECTRONIC DENSITY OF STATES



OTHER DECAYS (PLASMON, SCATTERING ON DEFECTS...)

UNITS

GAUSSIAN UNITS $[Q] = \text{STATC}$

COULOMB LAW $F = \frac{Q_1 Q_2}{r^2}$ $[\vec{E}] = [\vec{B}]$

ATOMIC UNITS

$m_e = 1$ $\hbar = 1$ $e = 1$ $\therefore \alpha = \text{FINE STRUCTURE CONSTANT} = \frac{e^2}{\hbar c} \approx \frac{1}{137}$

$$\Downarrow \\ c = 137$$

[ENERGY] = HARTREE = 2 RYDBERG = $2 \times (13.605 \text{ eV})$

[length] = $a_0 = \frac{\hbar^2}{m_e e^2} = \text{BOHR RADIUS} = 0.529177 \text{ \AA}$

HAMILTONIAN $\vec{B}=0$

$$H = \sum_i H_i^{1e} + \sum_{\substack{i,j \\ j \neq i}} H_{ij}^{2e}$$

$$H_i^{1e} = \frac{1}{2m_e} \vec{p}_i^2 + V(\vec{r}_i) + \frac{1}{2} \frac{1}{(m_e c)^2} \vec{\nabla} \cdot V(\vec{r}_i) \times \vec{p}_i \cdot \vec{S}_i$$

SPIN ORBIT

$$S_x |\uparrow\rangle = \frac{1}{2} \hbar |\uparrow\rangle$$

$$H_{ij}^{2e} = \frac{1}{2} \frac{e^2}{|\vec{r}_i - \vec{r}_j|} - \frac{1}{2} \frac{1}{(m_e c)^2} \frac{\vec{r}_i - \vec{r}_j}{|\vec{r}_i - \vec{r}_j|^3} \times \vec{p}_i \cdot \vec{S}_i$$

SPIN. ORBIT

SPIN-ORBIT HEURISTIC DERIVATION

$$\text{ELECTRIC FIELD } \vec{E} = + \frac{\vec{\nabla} V(\vec{r})}{e}$$

$$\text{ELECTRON VELOCITY} = \frac{\vec{p}}{m_e}$$

EFFECTIVE \vec{B} FIELD IN THE REFERENCE FRAME OF ELECTRON

$$\vec{B} = - \frac{\vec{v}}{c} \times \vec{E} = \frac{1}{m_e c} \vec{\nabla} V \times \vec{p}$$

$$H_H^{so} = g_s \frac{\mu_B}{\hbar} \vec{S} \cdot \vec{B} \quad g_s \approx 2 \quad \mu_B = \frac{e \hbar}{2 m_e c}$$

$$H^{so} = \left(\frac{1}{2} \right) H_H^{so} = g_s \frac{1}{2} \frac{e \hbar}{m_e c} \frac{1}{2 m_e c} = g_s \frac{1}{4 m_e^2 c^2} \vec{\nabla} V \times \vec{p} \cdot \vec{S}$$

→ THOMAS PRECESSION

$$g_s \approx 2 \quad H^{so} = \frac{1}{2 m_e^2 c^2} (\vec{\nabla} V \times \vec{p}) \cdot \vec{S}$$

• NOTATIONS

• ONE ELECTRON

DIRAC

$$|\overline{a}\rangle \in \left\{ \begin{array}{l} \text{SPACIAL \& SPIN} \\ \text{SPACE} \quad \text{SPACE} \end{array} \right\} \quad \text{FOR BREVITY}$$

$$|\overline{r}, \uparrow\rangle = |\vec{r}\rangle \otimes |\uparrow\rangle = |\vec{r}, \uparrow\rangle \quad \vec{r} \in \mathbb{R}^3$$

$$s = \uparrow, \downarrow = 1, -1$$

$$S_z |\uparrow\rangle = \frac{\hbar}{2} |\uparrow\rangle$$

SCHRÖDINGER

$$\psi(\vec{r}, s) \equiv \langle \vec{r}, s | \overline{A} \rangle$$

$$\text{OR } \psi_s(\vec{r}) \equiv \psi(\vec{r}, s)$$

SPINOR NOTATION (2 COMPONENT VECTOR)
SCHRÖDINGER DIRAC

$$\begin{pmatrix} \psi(\vec{r}, \uparrow) \\ \psi(\vec{r}, \downarrow) \end{pmatrix} \equiv \overline{\psi}(\vec{r})$$

• SIMPLIFIED SITUATION: FACTORIZED SPIN

IN ABSENCE OF SPIN-ORBIT

$$[S_z, H] = 0$$

\downarrow
#- SPIN-FACTORIZED

EIGEN-STATES OF H

$$|\overline{a}\rangle = |a\rangle \otimes |s\rangle \quad \text{IN SPINOR EITHER } \begin{pmatrix} \psi(r) \\ 0 \end{pmatrix} \text{ OR } \begin{pmatrix} 0 \\ \psi(r) \end{pmatrix}$$

NOT TRUE IN PRESENCE OF SPIN-ORBIT

OR ...

WITHIN A MEAN-FIELD APPROACH (SEE LATER)
INDIVIDUAL ELECTRONS IN THE SYSTEM-GROUND STATE
CAN BE SPINOR INSTEAD OF SPIN-FACTORIZED
EVEN IN ABSENCE OF SPIN-ORBIT

IN PARTICULAR - NON-MAGNETIC-N-ELECTRONS
FOR ANY AXIS: SPIN FACTORIZED

- MAGNETIC SYSTEMS

- COLINEAR FERROMAGNETISM
ANTI FERROMAGNETISM

SPIN-FACTORIZED
WITH $\vec{L} \parallel$ MAGNETISATION

- NON-COLINEAR MAGNETISM

SPIN-CAN NOT BE FACTORIZED
ELECTRON SHOULD BE REPRESENTED
WITH SPINORS

KRAMER'S DOUBLET

EVEN IN PRESENCE OF SPIN ORBIT
THE HAMILTONIAN IS INVARIANT UPON
TIME REVERSAL

T = TIME REVERSAL OPERATOR

$$T^2 = -1 \quad T^{-1} = -T$$

$$T^{-1} \vec{r} T = \vec{r} \quad T^{-1} \vec{p} T = -\vec{p} \quad T^{-1} \vec{S} T = -\vec{S}$$

IN SCHRÖDINGER REPRESENTATION

$$T = -i\sigma_y \underset{\substack{\downarrow \\ \text{complex} \\ \text{conjugate}}}{CC} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} CC$$

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

\Downarrow

$$T^{-1} H T = H$$

$$H^{1e} |\bar{a}\rangle = E_a |\bar{a}\rangle$$

$$T^{-1} H^{1e} T |\bar{a}\rangle = E_a |\bar{a}\rangle$$

$$H^{1e} T |\bar{a}\rangle = E_a T |\bar{a}\rangle$$

$$H^{1e} |\bar{a}'\rangle = E_a |\bar{a}'\rangle \quad |\bar{a}'\rangle = T |\bar{a}\rangle$$

$\{|\bar{a}\rangle, |\bar{a}'\rangle\}$ DEGENERATE KRAMER'S DOUBLET WITH
EIGENENERGY E_a

NOTE THAT $|a'\rangle$ AND $|a\rangle$ DISTINCT STATES

SINCE ${}^{\rightarrow}\langle \vec{r} | a \rangle = \begin{pmatrix} \psi(\vec{r}, \uparrow) \\ \psi(\vec{r}, \downarrow) \end{pmatrix}$

$$\begin{aligned} \overline{\langle a | a' \rangle} &= \overline{\langle a | T | a' \rangle} = \int d^3r \left(-\psi^*(\vec{r}, \uparrow) \psi^*(\vec{r}, \downarrow) + \psi^*(\vec{r}, \downarrow) \psi^*(\vec{r}, \uparrow) \right) \\ &= 0 \end{aligned}$$